

## Could a Massless SU(5) Theory Underly the Standard

**Model ?** { *Big X-Section New LHC Physics is Predicted !* } \*

Using supercritical RFT, I have argued that the uniquely unitary\* **Critical  $\mathbb{P}$**  occurs { via anomaly dynamics } in “**QCD<sub>S</sub>**” = **QCD<sub>*n<sub>f</sub>=6*</sub>** + { a sextet quark doublet } →

**EW sym-breaking** { sextet pions } &  
**Dark Matter** { stable sextet neutrons }

( Anomaly color factors  
=> big  $\mathbb{P}$  x-sections  
for EW bosons & sextet  
nucleons – at the LHC ? )

Remarkably, **QCD<sub>S</sub>** embeds uniquely in “**QUD**”\* - a massless SU(5) theory that, even more remarkably, may have (via massless fermion IR anomalies) a bound-state **S-Matrix** that reproduces the full **Standard Model**.

\* The RFT Critical  $\mathbb{P}$  (alone) satisfies all high-energy unitarity constraints.

\* **Quantum Uno/Unification/Unique/Unitary/Underlying Dynamics**

\* Presented at Gribov-80, May 2010, Trieste, Italy

**QUD**  $\longleftrightarrow$  **SU(5)** gauge theory with **massless, left-handed, fermions** in the  **$5 \oplus 15 \oplus 40 \oplus 45^*$**  representation.

**Uniquely** discovered as **1) anomaly free 2) asymptotically free {just} & 3) contains the EW symmetry-breaking sextet sector.** Under  $SU(3) \otimes SU(2) \otimes U(1)$

$$5 = (3, 1, -\frac{1}{3}) \{3\} + (1, 2, \frac{1}{2}) \{2\} , \quad 15 = (1, 3, 1) + (3, 2, \frac{1}{6}) \{1\} + (6, 1, -\frac{2}{3}) ,$$

$$40 = (1, 2, -\frac{3}{2}) \{3\} + (3, 2, \frac{1}{6}) \{2\} + (3^*, 1, -\frac{2}{3}) + (3^*, 3, -\frac{2}{3}) + (6^*, 2, \frac{1}{6}) + (8, 1, 1) ,$$

$$45^* = (1, 2, -\frac{1}{2}) \{1\} + (3^*, 1, \frac{1}{3}) + (3^*, 3, \frac{1}{3}) + (3, 1, -\frac{4}{3}) + (3, 2, \frac{7}{6}) \{3\} + (6, 1, \frac{1}{3}) + (8, 2, -\frac{1}{2})$$

Not only does **QUD** contain **QCD<sub>S</sub>**, both the triplet quark & lepton sectors {not asked for} are amazingly close to the SM !!! There are three “generations” – {1}, {2}, {3}.

Very importantly, **QUD is real {vector-like} wrt SU(3)xU(1)<sub>em</sub>.** **SU(2)xU(1) is not quite right but the lepton anomaly is correct**  $\Rightarrow$  to be physically realistic,

**all elementary leptons & quarks must be confined & massless !!**

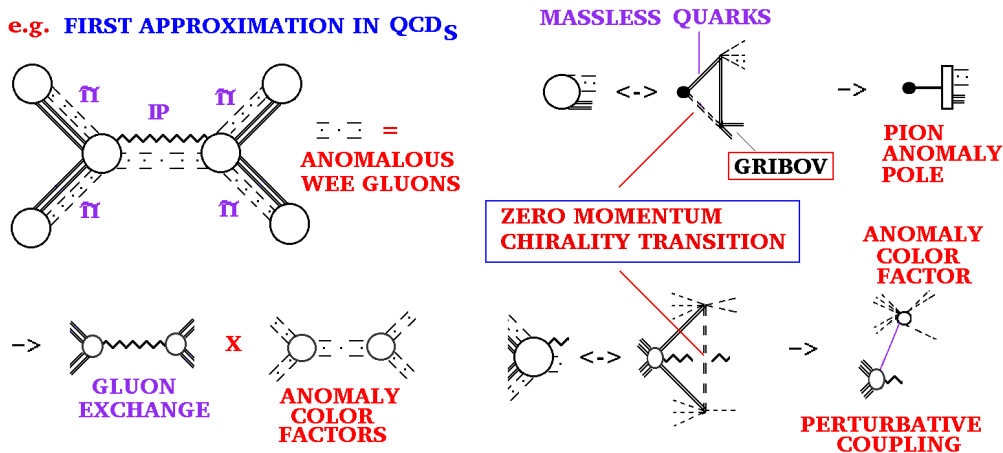
As for **QCD<sub>S</sub>**, access to the QUD bound-state S-Matrix is provided by **multi-regge theory !!!**

The multi-regge region involves multiple  $\infty$ -momenta that {fundamentally} allow wee partons to (simultaneously) play a vacuum role in both states & interactions.

I will build up the multi-regge QUD S-Matrix via IR divergences, by removing masses, & a  $k_{\perp}$  cut-off, in perturbative reggeon diagrams.

In massless\* QCD<sub>S</sub> - Anomalies + IR fixed pt.  $\rightarrow$  IR divergence  $\rightarrow$  Wee gluons in all bound-states & interactions - coupled via zero-momentum quark chirality transitions,

e.g. FIRST APPROXIMATION IN QCD<sub>S</sub>



- $\rightarrow$  Confinement & chiral sym brkg + parton model {in coexistence !!}
- $\rightarrow$  Quark hadrons  $\leftrightarrow$  no glueballs
- + additive quark model for  $\sigma_{\text{tot}}$ .
- $\rightarrow$  large sextet v. triplet anomaly color factor
- $\rightarrow$  Critical  $\mathbb{P}$  {regge pole + PPP}

\* Many massless mesons  $\rightarrow$  S-Matrix {???}. Masses require QUD.

**S-Matrix IR chirality transitions** *play an even more fundamental role in QUD. (They are produced by the zero fermion mass limit of reggeon anomaly vertices.)*

*Although only an outline, for which {as will be obvious} much further development is needed and many details are missing, my construction will imply the following.*

1. **All elementary fermions are confined.** *Bound-states are formed by anomaly poles*

**the symmetry breaking  $SU(5) \rightarrow SU(3) \otimes U(1)_{em}$**   
*is due to zero-momentum chirality transitions.*

2. **Interactions are vector boson reggeons + anomalous wee gauge bosons.**

3. **Symmetry-breaking is an S-Matrix phenomenon  $\leftrightarrow$  no off-shell amplitudes (?)**

*$SU(5)$  is unbroken at large  $k_{\perp}$  but, although QUD lies in the “conformal window”,*  
**the S-Matrix has only SM interactions & a spectrum of SM form.**

4. **All particles (including neutrinos) are bound-states with dynamical masses.**

**There is no Higgs !!**

**Because of it's uniqueness, QUD is either right or wrong - in it's entirety.**  
**{To be right} it must reproduce the full Standard Model S-Matrix !!**

{ Motivated by a unitary  $\mathbb{P}$  !! } QUD could provide a remarkably economic unification & even {perhaps} an origin\* for the SM. Beyond the SM generations, there is **only**

1. A sextet quark sector that minimally, & naturally, solves two major mysteries  
**Sextet pions**  $\rightarrow W^\pm, Z^0$  masses, & **stable\*** **sextet neutrons**  $\rightarrow$  **dark matter !**
2. A “lepton-like” **octet quark sector** that is buried in all states as UV anomaly poles  
 $\rightarrow$  **leptons & hadrons in SM generations**
3. A pair of exotically charged quarks.

**Nothing else !!**

Although the physics is both novel & radical, the outcome is simple, consistent with established SM physics, & explains many puzzles. But, the multi-regge theory that I use to uncover it is so erudite that general interest will probably require (what would surely be)

**A MAJOR EXPERIMENTAL  
DISCOVERY**

LHC  
please  
discover  
asap !!

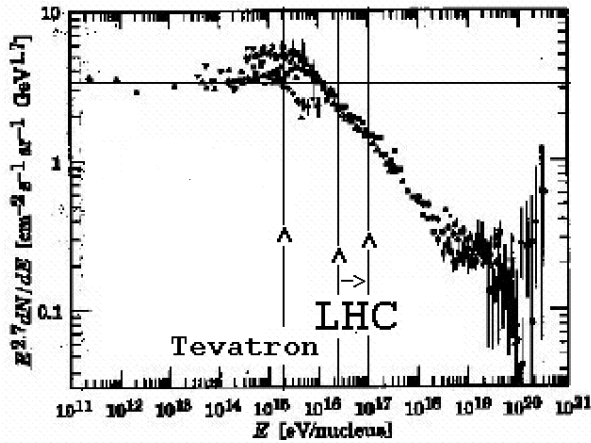
**BIG** x-sections for  
multiple  $Z'$ 's &  $W'$ 's,  $N_6$  &  $P_6$   
pairs, & **distinctively**,  $\gamma \mathbb{P} \rightarrow Z$   
**double  $\mathbb{P} \rightarrow ZZ, WW$  pairs**

\* Possibly, the SM (within QUD) is **THE** unique unitary particle S-Matrix !!

\* Sextet protons are unstable – decaying to sextet neutrons.

**COSMIC RAYS** already suggest that new large x-section physics including

**dark matter could appear at the LHC !!**



*The spectrum knee occurs between Tevatron and LHC energies. It is remarkably well-established, yet not understood. Although dark matter was unknown, a major threshold for neutral particles, unobserved in detectors, was initially suggested { ~ 40 years ago ! }. Underestimation of the energy would pile-up events as a “knee”.*

**If the dark matter x-section is large at the LHC, a link to the knee is surely inevitable !!**

For the sextet sector **three effects** should combine to produce a knee.

1. Prolific production of EW bosons increases  $\langle p_{\perp} \rangle$  dramatically (& increases neutrino production) – leading to energy underestimation.
2. **Threshold production of sextet neutron** dark matter {  $\leftrightarrow$  inclusive P }.
3. Sextet neutrons as ( UHE? ) incoming cosmic rays with a (P) threshold for atmospheric interaction not far below the knee.

X-sections must be **BIG** & at UHE must dominate  $\sigma_{tot}$  {  $\Rightarrow$  Dark Matter dominates }  
 early universe x-sections ?? }

**Related  
{crucial}  
questions.**

- {** i) Can **MASSLESS** QUD produce SM scales?  
ii) Why should the sextet sector have **BIG** x-sections? **}**

**An IR fixed-pt**  $\implies \alpha_{QUD}$  **is very small**  $\lesssim \frac{1}{120}$  (QUD is nearly conformal !)

$\implies$  **SM couplings**  $\longleftrightarrow$  **QUD evolution**  $\{\equiv \text{no off-shell amplitudes}\}$

But, {most likely}  $\implies$  **very small mass neutrinos** (no color/electric charge).

**All particles are bound-states**  $\implies \alpha_{QUD}$  has no physical meaning.

*S-Matrix amplitudes are selected by an IR divergence  $\longrightarrow$*

- *Physical states & amplitudes all contain infinite sums of wee gauge bosons involving anomaly color factors* {expressed, presumably, as integral formulae}.  
 $\implies$  *All interaction strengths are enhanced, with the SU(3) interaction strongly amplified by both color factors & the triple  $\mathbb{P}$  interaction* {see later}.
- *$\mathbb{P}$  anomaly color factors imply high-energy sextet x-sections are much larger than triplet x-sections*  $\Rightarrow$  **Dark Matter dominates early universe x-sections ??**

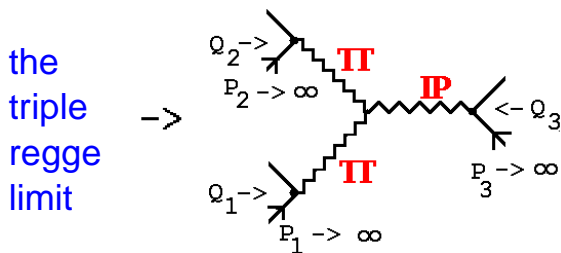
My analysis of QUD anomaly dynamics  
relies fundamentally on multi-regge theory. }

To give an outline, I begin with  
some basics - a regge pole gives

$$A^+(s, t) \underset{s \rightarrow \infty}{\sim} f(t) \frac{s^{\alpha(t)}}{\sin \frac{\pi}{2} \alpha(t)} \implies \begin{array}{l} t\text{-channel bound-state poles} \\ (\text{at } \alpha(t)=0, 2, \dots) \text{ can be discovered} \end{array}$$

by calculating the  $s \rightarrow \infty$  ( $\infty$ -momentum limit) in the cross-channel.

In multi-regge limits multiple regge poles appear, e.g.



$P_1, P_2, P_3 \rightarrow \infty$  along distinct light-cones.

In the  $P_3$  rest-frame, the bound-state regge pole pions also have  $\infty$ -momentum.

Continuation to  $Q_1^2 = Q_2^2 = m_\pi^2$  gives the on-shell pion amplitude for  $\mathbb{P}$  exchange.

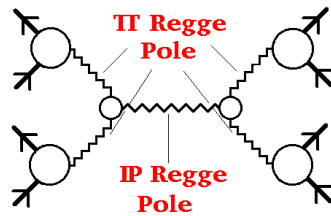
A priori, at  $\infty$ -momentum, **wee partons** could play a **vacuum** role\* in bound-states - if they are “universal”. We will see that indeed, in QUD, multi-regge reggeon diagrams do produce “vacuum wee partons”, but we must first introduce further  $\infty$ -momenta !!

\* c.f. light-cone quantization using the perturbative vacuum.



In the “**di-triple regge**” (DTR) limit two triple-regge limits are separated by a further  $\infty$ -momentum. Now, regge-pole  $\pi'$ s can scatter via the  $\mathbb{P}$ . All the  $\pi'$ s & the  $\mathbb{P}$  have  $\infty$ -momentum in some frame

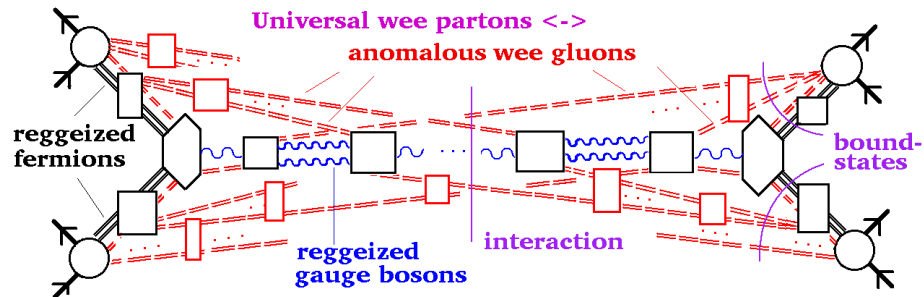
Hadron  
DTR  
Amplitude



$\Rightarrow$  both bound-states ( $\pi'$ s) & interactions ( $\mathbb{P}$ ) could appear as parton (reggeon) states if universal wee reggeons carry vacuum properties  $\{ \gg \text{parton model} \}$ . We will see

that, in QUD, an initial divergence produces “**anomalous wee gluons**” universally !!

Multiple  
wee parton  
interactions  
in a typical  
initial QUD  
DTR Amplitude



That SM states & interactions {& the **Critical P**} emerge from such diagrams as the complexity increases, is what has to be demonstrated !!

## Reggeon Diagrams {a crash course} -

1. In (multi-)regge limits the large light-cone momenta are routed through feynman diagrams so that internal particles are maximally close to mass-shell & have large relative rapidities.

→  $k_{\perp}$  diagram integral ( $\leftrightarrow$  close to on-shell lines contracted )  
multiplied by rapidity logarithms.

2. Summing rapidity logarithms via reggeon propagators → reggeon diagrams.  
{ Infinitely many feynman diagrams  $\longleftrightarrow$  1 reggeon diagram }
3. Internal particles with finite relative rapidity → couplings with more structure, including fermion loop interactions.
4. In a gauge theory triangle anomalies occur but {because a four-dimensional interaction is involved} only in special vertices coupling distinct reggeon channels - not\* in the single channel reggeon diagrams describing large  $k_{\perp}$  amplitudes.

\*SU(5) symmetry is not violated at “short distances” in QUD.

**Reggeon anomaly vertices** include axial-vector/vector/vector triangle diagrams  $T^{AVV}$  that, in QUD, must be defined at zero mass. At first sight,

- **chirality is conserved at zero mass**  $\implies T^{AAA} = T^{AVV} = T^{RRR} + T^{LLL}$   
 $\implies$  a conflict between the axial-vector anomaly & vector current conservation.

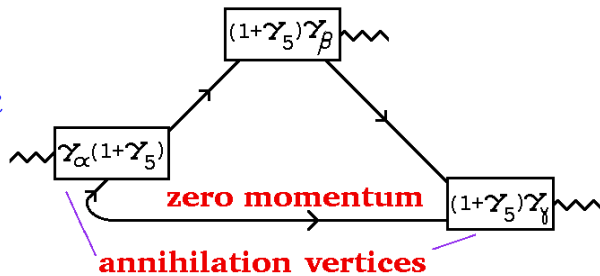
But, regularization of  $\gamma_5$  amplitudes is a major issue !!

- Fortunately, {it can be shown that} **vector current conservation + axial anomaly**  
 $\implies$  **unique massless anomaly pole amplitudes**  $\implies$

IR pseudoscalar anomaly poles appear similarly in both non chirality-violating & chirality-violating amplitudes via the triangle singularity.

e.g.

$T^{LRR}$



**Two annihilation vertices**  $\leftrightarrow$  **zero momentum chirality violation**  $\implies$  the pseudoscalar pole can be a chiral Goldstone boson.

**UV chirality transitions similarly produce anomaly poles as part of a Pauli–Villars subtraction when  $\lambda_\perp \rightarrow \infty$ .**

I use QUD reggeon diagrams to construct small  $k_{\perp}$  DTR amplitudes\*.

- I start with massive reggeons, via scalar VeV's, & with a  $k_{\perp}$  cut-off  $\lambda_{\perp}$  {VeV fermion masses, necessarily, identify particle/antiparticle pairs.}
- Because IR divergences produce wee partons in the massless theory, how the masses are removed is crucial. Anomalies play a key role {via the Gribov ambiguity}.

I take limits as follows -

1. Fermion masses  $\{24 \text{ \& } 5 \oplus 5^* \text{ VeV's}\}$  are removed first, leaving chirality transitions in anomaly vertices that conserve  $SU(3) \otimes U(1)_{em}$  but not  $SU(5)$ .
2. For gauge bosons,  $5 \oplus 5^* \text{ VeV's} \rightarrow$  smooth massless limit {complimentarity} Exponentiation of reggeization divergences confines\* global  $SU(5)$ , leaving IR finite interaction kernels. I take the limit in stages

$$\rightarrow SU(2)_C, \rightarrow SU(4), \lambda_{\perp} \rightarrow \infty, \rightarrow SU(5) \quad \left( \begin{array}{l} SU(2)_C \rightarrow SU(3)_C \\ \text{asymptotic} \\ \text{freedom} \end{array} \right) \quad \left( \begin{array}{l} \text{Supercritical} \\ \rightarrow \text{Critical P} \end{array} \right)$$

\* Very different from the large  $k_{\perp}$  use of QCD diagrams in BFKL physics.

\* Not true confinement! Multi-gluon singularities remain.

Because of  $\lambda_{\perp}$ , fermion loops do not have Ward identity zeroes when subsets of reggeon momenta vanish  $\Rightarrow$  many new divergences that exponentiate amplitudes to zero {including most left-handed boson couplings}.

**The first color restoration**  $\leftrightarrow$  (vector)  $SU(2)_C \longrightarrow$  a wee gluon divergence that couples only via anomaly vertices & so does not exponentiate  $\longleftrightarrow$

“anomalous wee gluons”  $\left\{ \begin{array}{l} \text{Sets of massless reggeized gluons, with all } k_{\perp} \text{'s scaled to zero} \\ \text{- with } I = 0 \text{ \& "anomalous" color parity } C \neq \tau = \text{signature.} \end{array} \right\}$

For  $SU(2)$ , only  $\tau = -C = -1$  is possible  $\leftrightarrow 3, 5, \dots \infty$  gluon reggeons.

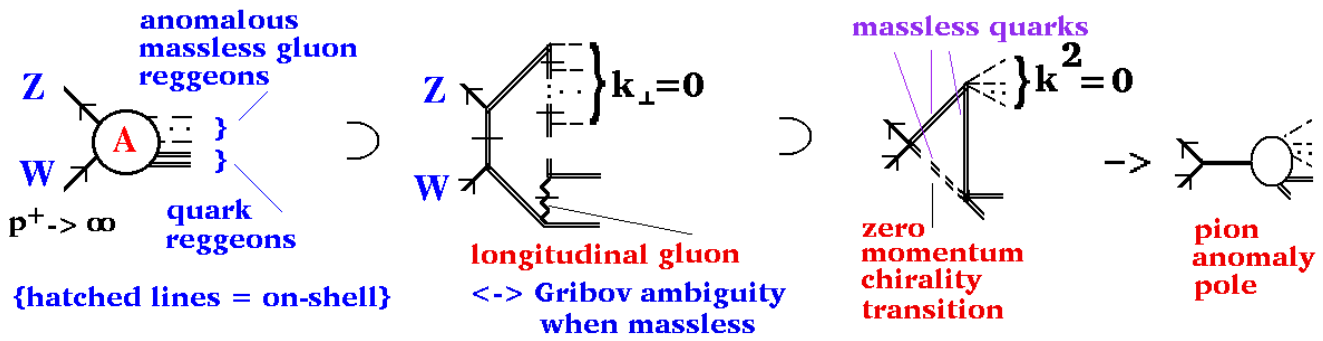
**IR fixed-point scaling**  $\implies$  iteration of  $I = 0$  reggeon kernels reproduces the basic divergence with a

$$\left( \sum \left\langle \begin{array}{c} \vdots \\ \vdots \end{array} \right| T_N \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| K_N^0 \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| T_N \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| K_N^0 \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| \right) T_N \left( \sum \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| K_N^0 \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| T_N \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| K_N^0 \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| T_N \left| \begin{array}{c} \vdots \\ \vdots \end{array} \right| \right)$$

factorized residue.  
Factorizing off the divergence

$\longrightarrow$  universal wee gluon component of reggeon states & interactions.

The divergence also produces **bound-state anomaly poles** in vertices, e.g.



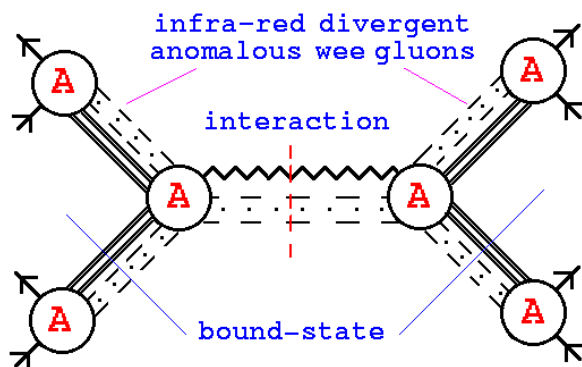
It is vital for S-Matrix symmetry breaking *that an anomaly pole connects*

1. a physical fermion pair + wee gluon “vacuum component”  $\leftrightarrow$  Reggeon state with color zero projection  $\rightarrow$  finite perturbative interactions, & {via the interaction}
2. an opposite chirality fermion pair - one unphysical & with zero momentum.  $\leftrightarrow$  pseudoscalar Goldstone boson\* associated with symmetry breaking.

**Effectively**, an anomaly pole bound–state is created by a zero momentum shift of the Dirac sea. **By absorbing anomalous wee gluons, a physical fermion makes a symmetry–breaking chirality transition to an unphysical “hole state”.**

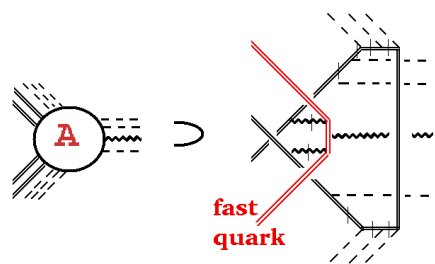
\* An  $\infty$ –momentum anomaly pole has physical Goldstone boson couplings.

The simplest  
DTR amplitudes  
selected by the  
anomalous wee  
gluon divergence.



=== = fermions  
 ~~~ = gauge boson

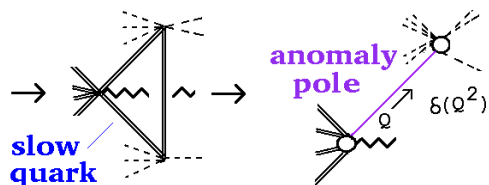
(A) = chirality transition anomaly pole vertex



**WEE GLUON ANOMALY COLOR FACTOR**  $\leftrightarrow$  sum over all wee gluons coupling to the slow quark loop

anomaly pole factorization

$\Rightarrow$  in QUD the P will couple more strongly to sextet states than to triplet states



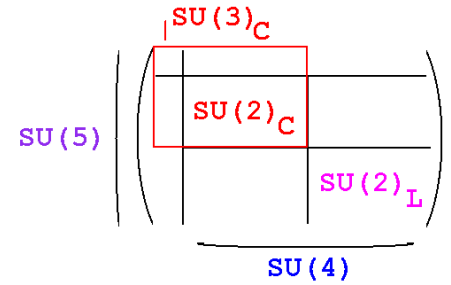
**PERTURBATIVE  $\leftrightarrow$  QCD ADDITIVE QUARK MODEL**

**Restoration of  $SU(2)_C$**  - gives anomaly pole chiral Goldstones ( $\pi_C$ 's) due to  $5 \oplus 5^*$  chirality transitions, that are  $q\bar{q}$  ("mesons") or  $qq$  &  $\bar{q}\bar{q}$  ("nucleons").

The  $q$ 's are **3's, 6's, & 8's** under  $SU(3)_C$ . **8's** are real wrt  $SU(3)_C$ , but contain complex chiral doublets wrt  $SU(2)_C$ . Via chirality transitions, the  $\pi_C$ 's are also reggeon states

$$\pi_C = q_L \bar{q}_R - \bar{q}_L q_R \rightarrow \begin{array}{c} \text{=====} \\ \text{-----} \end{array} \} q_L \bar{q}_L \text{ pair} + \dots$$

- Other reggeon states containing a  $\pi_C$  are also selected { will give leptons &  $SU(5)$  symmetry }



$$\left. \begin{array}{l} \text{=====} \} q_L \bar{q}_L \text{ pair} \\ \text{-----} \} \text{anomalous wee gluons} \\ \text{~~~~~} \} \text{massive gauge boson reggeons} \\ \text{=====} \} \text{fermion reggeons} \end{array} \right\} \leftrightarrow \pi_C$$

$SU(2)_C$   
color zero

To avoid fermion loop exponentiation of the anomaly divergence the massive gauge boson reggeons must be vectors  $\leftrightarrow SU(2)_C$  singlet gluons or photons.

**Interactions** are even signature & are  $SU(2)_C$  singlet massive vector exchange, together with anomalous wee gluons  
 $\longrightarrow$

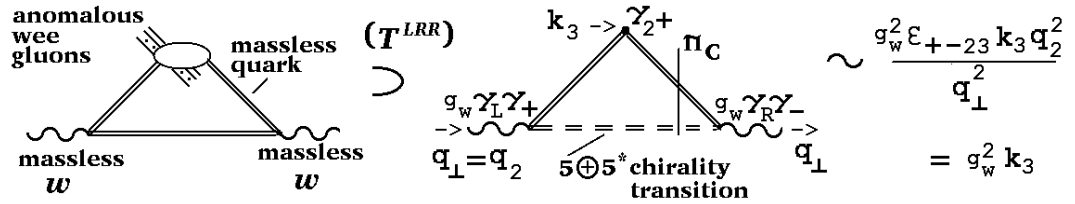
$$\begin{array}{l} \text{~~~~~} \} \text{massive gauge boson reggeon} \\ \text{-----} \} \text{anomalous wee gluons} \end{array}$$



- Massive  $SU(3)_C$  gluon exchange in the  $SU(2)_C$  wee gluon background  $\leftrightarrow \mathbb{P}$ .
- The massive gluon can be replaced by a massive  $\gamma$ ,  $W^\pm$  or  $Z^0$ .

Elementary left-handed exchanges ( $W^\pm$  &  $Z^0$ ) exponentiate to zero, but

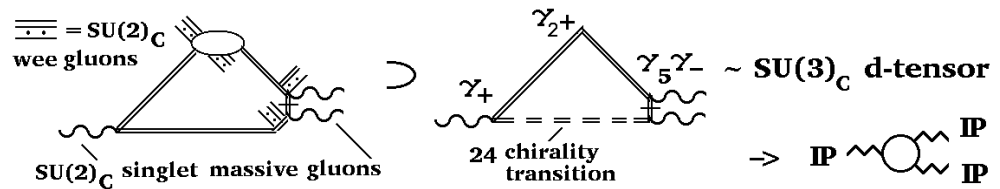
$5 \oplus 5^*$  chirality  
transitions  
 $\rightarrow$  crucial  
wee-gluon  
interactions.



- $\rightarrow$  mass =  $M_W^2 \sim g_W^2 \int dk k \leftrightarrow$  universal wee gluon integral.
- $\leftrightarrow$  mixing with the  $\pi_C'$ s {  $\pi_6$  dominates after  $SU(3)_C$  is restored }
- $\rightarrow$  new quantum number {  $\rightarrow$  sextet flavor } & no exponentiation to zero.
- $\rightarrow$  massive vector  $W^\pm$  &  $Z^0$

Odd-signature interactions are  $\gamma \mathbb{P}$ ,  $W^\pm \mathbb{P}$  &  $Z^0 \mathbb{P}$  {  $\rightarrow$  physical  $\gamma$ ,  $W^\pm$ ,  $Z^0$  after  $SU(3)_C$  restoration }.

24 chirality  
transitions  
→ d-tensor  
triple  $\mathbb{P}$  vertex.



- Wee gluons → orthogonal  $\gamma$ -matrices  $\leftrightarrow \gamma_5$ .  
→ essential vertex for the SU(3)<sub>C</sub> interaction & the Critical  $\mathbb{P}$ .

**SU(4) Restoration**  $\longleftrightarrow$  SU(2)<sub>C</sub> singlet vector  $\gamma_4$  becomes massless

- Other gauge bosons (that become massless) have left-handed couplings & appear only in interaction kernels {reggeon diagrams are exponentiated to zero.}
- $\gamma_4$  pairs have 1) even signature, 2) an SU(4) singlet projection, & 3) exponentiate to zero all amplitudes except when coupling to anomaly poles.  $\{\leftrightarrow 1 - e^{-\infty}\}$
- Anomalous wee bosons  $\{\gamma_4$  pairs + SU(2)<sub>C</sub> anomalous wee gluons}  
→ even signature  $\gamma$  - with SU(4) singlet projection.
- Anomalous wee bosons + massive gluon → Supercritical  $\mathbb{P}$ .

**Pseudoscalar anomaly poles** coupling to  $\gamma_4$  pairs are produced by

1. *lepton pairs*  $(1, 2, \frac{1}{2})$  &  $(1, 2, -\frac{1}{2})$ .

Chiral symmetry + **24** chirality transitions  $\rightarrow$  **pseudoscalars**  $\pi_L^{\pm,0}$

2.  $SU(2)_C$  *singlets*  $(8, 1, 1)$  &  $(8, 2, -\frac{1}{2})$ .

Chiral symmetry +  $5 \oplus 5^*$  chirality transitions  $\rightarrow$  **pseudoscalars**  $\eta_8^{\pm,0}$ .

Now, the anomalous wee boson divergence  $\implies$  **physical bound-states**

1. **contain two pseudoscalar anomaly poles** coupling to the divergence
2. **have  $SU(4)$  singlet projections** as reggeon states.
3. **Fermion bound-states** contain an additional elementary fermion.

**Leptons** -  $\pi_L + \pi_8 + \text{additional lepton} \rightarrow 3$  lepton generations.

**Hadrons** -  $\pi_{3,6} + \eta_8 \rightarrow \text{mesons}, + \text{additional quark} \rightarrow \text{baryons}.$

{ More details after  $SU(5)$  restoration. }

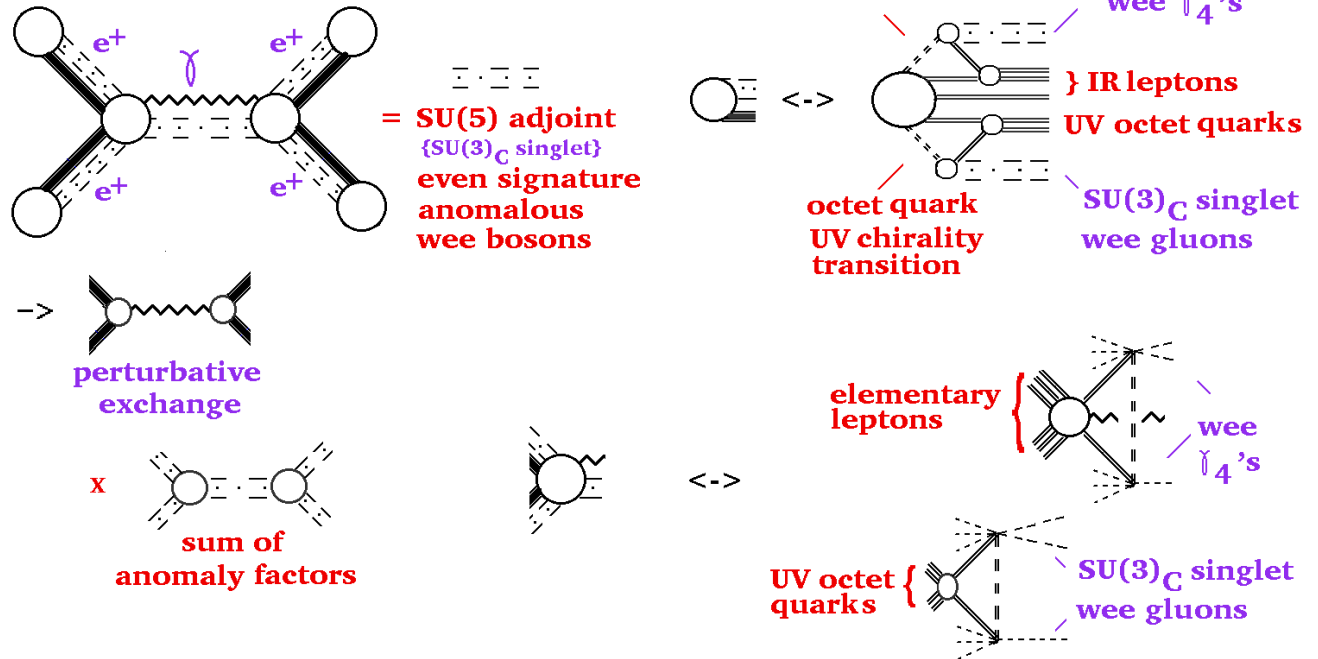
## SU(5) Restoration $\leftrightarrow$ *SU(4) singlet vector becomes massless.*

1. **Supercritical  $\mathbb{P}$**   $\{\leftrightarrow$  SU(4) singlet + odd-signature anomalous wee bosons with SU(3)<sub>C</sub> color  $\} \longrightarrow$  **Critical  $\mathbb{P}$**   $\{\leftrightarrow$  randomizing symmetry breaking  $\}$
2. **Even signature anomalous wee bosons** {with SU(3)<sub>C</sub> color zero}  $\longrightarrow$  **odd-signature massless  $\gamma$  & massive  $W^\pm, Z^0$ .**
3.  **$\pi_8$  &  $\eta_8$  combine in real octet SU(3)<sub>C</sub> representation ( $\Pi_8$ )  $\implies$** 
  - **octet IR anomaly poles cancel in all amplitudes,**
  - **UV  $\Pi_8$  poles appear as  $\lambda_\perp \rightarrow \infty$   $\{\leftrightarrow$  Pauli-Villars subtractions  $\}$** $\implies$  **Leptons & hadrons have IR components combined with UV  $\Pi_8$  's.**
4. **Interactions are determined by IR components - with states expressed as physical reggeons with an SU(5) singlet projection.**
5. **UV  $\Pi_8$  anomaly poles  $\implies$  IR fermions must be**

$$(2, -\frac{1}{2})_L \text{ or } (2, \frac{1}{2})_R \text{ or } (1, 1)_L \text{ or } (1, -1)_R \quad \{SU(2) \otimes U(1)\}$$

$$\implies$$
 **IR components of leptons & hadrons form SM generations.**

## First approximation in QUD



**Octet quarks coupled via SU(3) $_C$  color zero anomalous wee gluons**  
 {contained in SU(5) adjoint even–signature anomalous wee bosons}  
**provide a UV light-cone contribution in all states & vertices.**

The  $SU(3) \times SU(2)_L \times U(1)$  lepton content, expressed as multi-fermion states, is

- $(e^-, \nu) \leftrightarrow (1, 2, -\frac{1}{2}) \times \pi_L^0 \times \Pi_8$   
 $\leftrightarrow (1, 2, -\frac{1}{2}) \times (1, 2, -\frac{1}{2}) (1, 2, \frac{1}{2}) \times (8, 1, 1) (8, 2, -\frac{1}{2})$   
 $\leftrightarrow \text{SU}(5) \text{ singlet/adjoint} - 45^* \times 45^* \times 5 \times 40 \times 45^*$
- $(\mu^-, \nu) \leftrightarrow (1, 2, \frac{1}{2}) \times \pi_L^- \times \Pi_8$   
 $\leftrightarrow (1, 2, \frac{1}{2}) \times (1, 2, -\frac{1}{2}) (1, 2, -\frac{1}{2}) \times (8, 1, 1) (8, 2, -\frac{1}{2})$   
 $\leftrightarrow \text{SU}(5) \text{ singlet/adjoint} - 5 \times 45^* \times 45^* \times 40 \times 45^*$
- $(\tau^-, \nu) \leftrightarrow (1, 2, -\frac{3}{2}) \times \pi_L^+ \times \Pi_8$   
 $\leftrightarrow (1, 2, -\frac{3}{2}) \times (1, 2, \frac{1}{2}) (1, 2, \frac{1}{2}) \times (8, 1, 1) (8, 2, -\frac{1}{2})$   
 $\leftrightarrow \text{SU}(5) \text{ singlet/adjoint} - 40 \times 5 \times 5 \times 40 \times 45^*$

$e^+, \mu^+, \tau^+$  can be obtained via charge conjugation.

## QCD within QUD - the states are

1. triplet mesons & baryons
2. sextet “pions” & “nucleons” ( $P_6$  &  $N_6$ )
3. no hybrid sextet/triplet quark states
4. no glueballs.

Consistent with, but fewer states than just confinement & chiral symmetry breaking.

- Sextet pions  $\rightarrow W^\pm$  &  $Z^0 \implies$  sextet nucleons are the **only new states**.
- Sextet quarks have zero current mass  $\implies N_6$  **is stable**  $\rightarrow$  **DARK MATTER**  
{electric charge  $\implies P_6$  is heavier - in contrast to the triplet sector}
- Critical  $\mathbb{P}$  & parton model,  $\longleftrightarrow$  **no BFKL pomeron & no odderon**.

Compared to conventional QCD, the states are fewer & the interaction is simpler - in agreement with experiment !!

$N_6$ 's have a very strong, very short range, QCD self-interaction  $\{\rightarrow$  clumps? $\}$  & a QCD interaction with normal matter only at UHE.  $N_6$  production will dominate UHE x-sections & early universe stable matter formation (& explain the CR knee!)

**The bound-state mass spectrum** will involve a combination of effects.

1. *Perturbative reggeization is a very small effect, since  $\alpha_{QUD}$  is so small.*
2. **Wee gluon anomaly interactions** {which need to be catalogued} **will mix the reggeon states & introduce anomaly color factors.**
3. The SU(3) strong interaction will be emphasized by both **color factors & by the resulting high mass sector.** Electroweak charges will also play a role.
4. **There is no symmetry that would conflict with the SM mass spectrum.**

A wide range of scales should emerge & in bound-states **all fermions will, effectively, have constituent masses.** *Connecting the  $\eta_6$  to top production suggests*

$$m_{q_6} \sim m_{top} \implies m_{N_6} \sim 500 \text{ GeV}$$

**In general, a better understanding of anomaly interactions & related wee gluon distributions is needed to determine if, & how many, parameters are involved.**

**Presumably, CP violation can be introduced via the anomalies, but is it essential?**



## QUD Virtues {beyond QCD}

It would be hard to over-emphasize the scientific {& aesthetic} importance of an underlying massless field theory for the Standard Model. In addition,

1. *QUD is self-contained - with only SM Interactions !!*
2. *The only new physics is the strong interaction sextet sector - giving **EW symmetry breaking, dark matter**, & unification without supersymmetry !!*
3. *The parity properties of the strong & weak interaction are naturally explained.*
4. *Particles & fields are truly distinct. Hadrons & leptons have equal status. **No off-shell amplitudes & no Higgs field**  $\leftrightarrow$  all symmetries & masses are S-Matrix properties.*
5. *The QUD S-Matrix could be the only “non-perturbative” part of field theory needed - with infinite momentum physics retaining a “perturbative” diagrammatic description.*
6. *Anomaly interactions mix the reggeon states &, presumably, introduce parameters.*
7. *Anomaly color factors should produce a wide range of scales & masses that could produce the SM spectrum - there is no conflicting symmetry.*
8. **Small neutrino masses** *should be due to the underlying small coupling.*
9. *There is no proton decay, but the **SU(5) Weinberg angle should hold!***
10. *Perturbatively, QUD is an asymptotically free, massless, fixed-point theory that should have no perturbative renormalons  $\implies$  it has no vacuum energy & would induce Einstein gravity with zero cosmological constant.*